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INVERSE MEAN FREE PATH, STOPPING POWER, CSDA RANGE, AND STRAGGLING IN ALLMINUM AND ALUMINUM OXIDE FOR ELECTRONS OF ENERGY = OR < 10 keV

J. C. Ashley, et al

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Prepared for:

Air Force Cambridge Research Laboratories Energy Research and Development Administration

December 1975

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December 1975

Scientific Report No. 1

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This research was sponsored by the Defence Nuclear Agency under Subtack Z99CaXTAD40, Work Unit 51, entitled "Secondary Electron Transport Planennerslogy" aid in part by the Emergy Research and Development Administration under contract with Union Carbide Corp.

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REPORT DOCUMENTATION	PAGE	READ INSTRUCTIONS BEFORE COMPLETING FORM
AFCRL-TR-75-0583	2. GOVT ACCESSION MO	J. RECIPIENT'S CATALOG NUMBER
. TITLE (and Substite) NVERSE MEAN FREE PATH, STOPP	1	5. TYPE OF REPORT & PERIOD COVERED
SDA RANGE, AND STRAGGLING IN.	ALUMINUM AND	Scientific - Interim
ALUMINUM OXIDE FOR ELECTRONS	OF ENERGY	6. PERFORMING ORG. REPORT NUMBER
≤ 10 keV		Scientific Report No. 1
. AUTHOR(e)		B. CONTRACT OR GRANT NUMBER(s)
J. C. Ashley R. H. Ri C. J. Tung V. E. Anderson	tchie	¥75-886
PERFORMING ORGANIZATION NAME AND ADDRESS		10. PROGRAM ELEMENY, PROJECT, TASK AREA & WORK UNIT NUMBERS
Oak Ridge National Laboratory Oak Ridge, Tennessee 37830		62704H CDNA0008
1. CONTROLLING OFFICE NAME AND ADDRESS		12. REPORT DATE
Air Force Cambridge Research Labor	ratories	December 1975
Hanscom AFB, Massachusetts 01731		15. NUMBER OF PAGES
Contract Monitor: Dr. John N. Brad		54
4. MONITORING AGENCY NAME & ADDRESS/If differen	nt from Controlling Office)	15. SECURITY CLASS. (al this report)
		Unclassified
		15. DECLASSIFICATION/DOWNGRADING SCHEDULE

Approved for public release; distribution unlimited.

17. DISTRIBUTION STATEMENT (of the abstract entered in Block 20, If different from Report)

This research was sponsored by the Defense Nuclear Agency under Subtask Z99QAXTA040, Work Unit 51, entitled "Secondary Electron Transport Phenomenology and in part by the Energy Research and Development Administration under contract with Union Carbide Corp.

19 KEY WORDS (Continue on reverse side if necessary and identify by block number)

electron transport energy loss theory

20. ABSTRACT (Continue on reverse side if necessary and identify by block number)

The interaction of electrons with the solids Al and Al<sub>2</sub>O<sub>3</sub> is described based on the electron gas model for the conduction band electrons in Al, a model insulator theory for the valence electrons in Al<sub>2</sub>O<sub>3</sub>, and inner shell ionization derived from atomic, general zed oscillator strengths. Contributions to the inverse mean free path and stopping power from the various interaction processes are tabulated for electron energies from 0.5 eV to 10 keV for Al and from 10 eV to 10 keV for Al<sub>2</sub>O<sub>3</sub>.

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Block 20. (Continued)

Electron range in the continuous-slowing-down approximation and straggling are tabulated for electron energies from 10 eV to 10 keV for both materials.

SECURITY CLASSIFICATION OF THIS PAGE(When Data Entered)

# CONTENTS

I.	INTRODUCTION
n.	GENERAL FORMULATIONS
ш.	DIMFP'S (DIFFERENTIAL INVERSE MEAN FREE PATHS) FOR THE
	ALUMINUM CONDUCTION BAND
rv.	DIMFP'S FOR INNER SHELLS
v.	DIMFP'S FOR VALENCE ELECTRONS IN A1203
	EXCHANGE CORRECTED DIMFF'S AND FORMULAE FOR THE
	TABULATIONS
νII.	REFERENCES
vIII.	ALUMINUM: EXPLANATION OF TABLES
	TABLE 1A - INVERSE MEAN FREE PATH OF ELECTRONS IN
	ALUMINUM 22
	TABLE 1B - STOPPING POWER OF ALUMINUM FOR ELECTRONS 27
	TABLE 1C - CSDA RANGES AND STRAGGLING OF ELECTRONS
	IN ALUMINUM
ıx.	ALUMENUM OXIDE: EXPLANATION OF TABLES
	TABLE 2A - INVERSE MEAN FREE PATH OF ELECTRONS IN A1203 39
	TABLE 2B - STOPPING POWER OF Al <sub>2</sub> O <sub>3</sub> FOR ELECTRONS 44
	TABLE 2C - CSDA RANGE AND STRAGGLING OF ELECTRONS
	IN A1 <sub>2</sub> O <sub>3</sub>

### I. INTRODUCTION

A quantitative description of the interaction of electrons with matter over a large range of energies is a subject of basic importance in a wide variety of theoretical and applied areas. From the theoretical standpoint, calculations of energy loss and range of electrons in many different materials have formed the basis of at least two extensive tabulations. Both of these works are restricted to electron energies  $\geq 10$  keV and are based on the Bethe theory of stopping power including various modifications and corrections (e.g. density-effect corrections). We feel that similar tabulations for electron energies  $\leq 10$  keV, based on a priori calculations using currently available theoretical information, will provide useful guides for interpretation of experimental data as well as input for calculations in applied areas.

Our work here will involve model calculations to describe the extended electron states of a solid (valence band or conduction band). The more tightly bound, inner shells of the atoms in the solid will be assumed to be essentially unchanged in character from those in free atoms. Thus, excitation of electrons from the inner shells will be based on calculations of atomic, generalized oscillator strengths (GOS's). More specifically, the differential inverse mean free path (DIMFP), which forms the basic function required in our work, will be derived from an electron gas model describing the conduction band electrons in Al and from a model insulator theory applied to the valence band in Al<sub>2</sub>O<sub>3</sub>. The DIMFP's for electron interaction with the inner shells of the Al and O atoms

are calculated from a priori atomic GOS calculations for exciting electrons to the continuum from the 2s and 2p sub-shells of Al<sup>3</sup> and the K shell of O.<sup>4</sup> The DIMFP for removing an electron from the K shell of Al is obtained from GOS values based on hydrogenic wavefunctions.<sup>5</sup> Given the DIMFP's associated with the most important electron interaction processes in the solids we then calculate inverse mean free paths, stopping powers, csda ranges, and range and energy straggling for electron energies from a few eV to 10 keV.

The details of the components of our calculations are described more fully in the next four sections. In section VI exchange corrections are discussed, expressions given for the exchange corrected DIMFP's, and formulas used in these tabulations are displayed.

#### II. GENERAL FORMULATIONS

A charged particle passing through a solid interacts with a large number of electrons simultaneously and it is thus appropriate to speak of a mean free path of the charged particle against energy loss to the solid. Assuming the effect of the charged particle on the medium may be treated in first Born approximation, the inverse mean free path, differential in momentum transfer,  $h\vec{k}$ , and energy transfer,  $h\omega$ , for a particle of velocity v is given by

$$\frac{d^2\mu}{dkd\omega} = \frac{2e^2}{\pi\hbar v^2} \frac{1}{k} \operatorname{Im}\left[\frac{-1}{\epsilon(k,\omega)}\right]$$
 (1)

where  $\epsilon(k,\omega)$  is the exact dielectric function of the solid.<sup>6,7</sup> We assume in this work that the solid is isotropic and homogeneous.

For our calculations of inverse mean free path, stopping power, etc., it is sufficient to compute inverse mean free paths differential in energy transfer only. This differential inverse mean free path (DIMFP) for energy loss  $\hbar\omega$  by an electron with energy  $E = mv^2/2$  in the solid is given by

$$\tau(E, h\omega) = \frac{d\mu}{d(h\omega)} = \frac{1}{\pi_a E} \int_{k}^{k_+} \frac{dk}{k} \operatorname{Im} \left[ \frac{-1}{\epsilon(k, \omega)} \right]$$
 (2)

where  $hk_{\pm} = \sqrt{2m} \left[ \sqrt{E \pm \sqrt{E - h\omega}} \right]$  and  $a_0 = h^2/me^2$ . This expression assumes that the energy-momentum relation for a swift electron in the solid does not differ appreciably from that of a free electron in vacuum.

Given  $\mathfrak{C}(k,\omega)$  for the solid, the quantities of interest here follow directly from  $\tau(E,\hbar\omega)$ . The inverse mean free path of the electron,  $\mu$ , is given by integrating over allowed energy transfers as

$$\mu(E) = \int d(h\omega) \ \tau(E, h\omega) \ . \tag{3}$$

The rate of energy loss of the electron, or the stopping power of the medium, is given by

$$S(E) = -dE/dx = \int d(h\omega) h\omega \tau(E, h\omega)$$
,

and the mean square energy loss per unit path length by

$$\Omega^{2}(E) = \int d(h\omega)(h\omega)^{2} \tau(E, h\omega) . \qquad (5)$$

With these results we may calculate the range of an electron in the continuousslowing-down approximation (csda range) by

$$R_{o}(E) = \int_{E_{o}}^{E} dE'/S(E') . \qquad (6)$$

The lower limit on this integration will be discussed further in Section VI. The mean square fluctuation in the range or "range straggling" will be calculated from Eq. (3) and Eq. (4) as 8

$$(R - R_o)_{AV}^2 = \int_{E_o}^{E} dE' \Omega^2(E') / [S(E')]^3 . \qquad (7)$$

In practice, the DIMFP will be evaluated as a sum of contributions from various distinct processes. For example we calculate a DIMFP for removing an electron from a given inner shell, a DIMFP for plasmon excitation, etc. The total DIMFP used to describe the interaction of an electron with the given solid will be given by

$$\tau(E,\omega) = \sum_{i} \tau_{i}(E,\omega)$$
 (8)

where the sum over i adds the contributions from the various interaction processes. The evaluation of the  $\tau_1$ 's for Al and Al<sub>2</sub>O<sub>3</sub> is described in the next three sections.

# III. DIMFP'S FOR THE ALUMINUM CONDUCTION BAND

The conduction band of Al will be described by an electron gas model. The dielectric response function in the form given by Lindhard $^6$  is

$$\epsilon(k,\omega) = 1 + (\chi^2/z^2) [f_1(x,z) + if_2(x,z)]$$
 (9)

in terms of the dimensionless variables  $x = \hbar \omega / E_r$  and  $z = k/2k_r$  with  $E_r$  the Fermi energy and  $k_r$  the Fermi wavenumber of the electron gas. Also,  $\chi^2 = e^2/\pi \hbar v_r$  where  $v_r$  is the Fermi velocity. The function  $f_1$  is given by

$$f_{1}(x,z) = \frac{1}{2} + \frac{1}{8z} \left[ 1 - (z - x/4z)^{2} \right] \ln \left| \frac{z - x/4z + 1}{z - x/4z - 1} \right| + \frac{1}{8z} \left[ 1 - (z + x/4z)^{2} \right] \ln \left| \frac{z + x/4z + 1}{z + x/4z - 1} \right|.$$
 (10a)

The function f<sub>2</sub> is defined by

$$f_{2}(x,z) = \begin{cases} \pi x/8z, & \text{for } z + x/4z < 1 \\ \frac{\pi}{8z} \left[ 1 - (z - x/4z)^{2} \right], & \text{for } |z - x/4z| < 1 < z + x/4z \quad (10b) \\ 0, & \text{for } |z - x/4z| > 1 \end{cases}$$

From Eq. (2), the DIMFP is given by

$$\tau(E, h\omega) = \frac{\chi^2}{\pi a_0 E} \int_{z}^{z_+} dz \frac{zf_2}{(z^2 + \chi^2 f_1)^2 + (\chi^2 f_2)^2}$$
(11)

where  $z_{\pm} = k_{\pm}/2k_{\tau}$ . Energy losses of an incident electron to an electron gas may be divided into: (a) losses resulting from excitation of single electrons out of the Fermi sea and (b) losses to collective oscillations of the electron gas (plasmon excitation). The contribution to the total DIMFP due to (a) is identified with the region in which  $f_2$  is non-zero. This electron-electron term is written as

$$\frac{d\mu_{ee}}{dx} \equiv \tau_{ee}(\varepsilon, x) = \frac{\chi^2 \theta(\varepsilon - x)}{\pi_{a_0}(\varepsilon + 1)} \int_{\frac{1}{2}(\sqrt{x+1} - 1)}^{\frac{1}{2}(\sqrt{x+1} + 1)} \frac{zf_2}{(z^2 + \chi^2 f_1)^2 + (\chi^2 f_2)^2}$$
(12)

where  $\epsilon$  is the incident electron energy measured from the Fermi level in units of the Fermi energy, i.e.,  $\epsilon = (E - E_f)/E_f$ , and the step function  $\theta(\epsilon - x)$  restricts the incident electron to states above the Fermi level.

The plasmon contribution arises from integration of Eq. (11) in the region where  $f_2 \to 0$ . The integral is zero except on the plasma resonance line defined by  $z = z_0(x)$  through the equation

$$F(x,z_0) = z_0^2 + \chi^2 f_1(x,z_0) = 0.$$
 (13)

The DIMFP for plasmon excitation is found to be

$$\tau_{\text{pl}}(\varepsilon,x) = \frac{1}{a_0(\varepsilon+1)} \frac{z_0}{|dF/dz|} \theta(x-x_{\text{min}}) \theta(x_{\text{mex}}-x)$$
 (14)

where x and x are the solutions of the equations

$$x_{\min} = 4z_{o}(x_{\min}) \left[ \sqrt{\epsilon + 1} - z_{o}(x_{\min}) \right]$$
 (15)

and

$$x_{max} = 4z_{o}(x_{max})[1+z_{o}(x_{max})].$$
 (16)

For our later discussion of exchange corrections we need the DIMFP for creation of secondary electrons in this model. The results of Ref. (7) lead to

$$\tau_{s}(\epsilon, \epsilon') = \frac{\chi^{2}}{8a_{o}(\epsilon+1)} \int_{\epsilon'}^{\min(\epsilon'+1, \epsilon)} dx \int_{z}^{z} \frac{dz}{z^{4} |\epsilon(2k_{F}z, E_{F}x)|^{2}}$$
(17)

where min(a,b) represents the smaller of the quantities a and b.

The use of this electron gas model, though based implicitly on the assumption of high electron densities, gives unexpectedly good results for real metal conduction band densities. In addition, Eq. (1) is derived in first Born approximation and would be expected to fail when the velocity of the incident electron approaches that of electrons in the electron gas. Lindhard, however, points out that this equation may be reasonably good at any velocity since: (a) the relative velocity of the incident electron and representative electrons in the electron gas remains of the order of the Fermi velocity no matter how slowly the incident electron moves, and (b) at low velocities the Coulomb field of the incident electron is strongly screened by polarization in the electron gas so \*! It for many purposes it can be considered small in the sense of perturbation theory.

### IV. DIMFP'S FOR INNER SHELLS

From a general expression for the dielectric function of a homogeneous, isotropic system  $^9$  we may show for values of  $\omega$  which correspond to ionization of the  $i^{th}$  inner shell in a solid that

$$\operatorname{Im}\left[\frac{-1}{\epsilon(k,\omega)}\right] \approx \operatorname{Im} \epsilon(k,\omega) \approx \frac{2\pi n_i e^2}{m\omega} \frac{\mathrm{d}f_i(k,\omega)}{\mathrm{d}\omega} \tag{18}$$

where  $df_i/d\omega$  is the generalized oscillator strength (GOS) for transitions from the i<sup>th</sup> independent inner shell. Here n is the number of i<sup>th</sup> inner shells per unit volume in the given solid. Equation (2) thus leads to

$$\tau_{i}(E, h\omega) = \frac{8\pi a_{o}^{2} n_{i}}{(E/R)(h\omega/R)} \int_{k}^{k_{+}} \frac{dk}{k} \frac{df_{i}(k, \omega)}{d(h\omega)}$$
(19)

where  $\hbar \omega$  is the energy transfer and  $R = e^2/2a_0$ .

Generalized oscillator strengths for the ionization of L shell electrons in Al have been calculated by Manson susing a nonrelativistic Hartree-Slater central field model of the atom. These GOS values have been used as input for numerical evaluation of the integral over momentum transfer in Eq. (19) to obtain differential cross sections,  $\tau_i/n_i = d\sigma_i/d(\hbar\omega)$ , for the 2s and 2p subshells of Al. These cross sections when multiplied by the appropriate value of  $n_i$  give the DIMFP for the particular material. For example,  $\tau_{2s}$  would denote the DIMFP for removing a 2s electron from an Al atom with the density of Al atoms corresponding to that in Al<sub>2</sub>O<sub>3</sub>. This same type of notation will be used for contributions to the inverse mean free path and stopping power. The binding energy for 2s and 2p electrons in Al is taken to be 118.524 eV and 80.88 eV, respectively, for these calculations.

Similar calculations have been done for the 1s shell in oxygen. The values of the GOS for excitation to the continuum from the 1s shell of Owere taken from the work of McGuire. The binding energy of the 1s shell in O is taken to be 536.6 eV.

Our calculation of the DIMFP for excitation of electrons from the K shell of Al to the continuum is based on cross sections derived using hydrogenic wave functions. A detailed discussion of this type of cross section calcuation is

given in Ref. 5. We have used the equations in this reference to calculate the DIMFP based on a binding energy for the K shell in Al of 1545.8 eV and an "effective atomic number" of  $Z_a = 12.5643$ .

# V. DIMFP'S FOR VALENCE ELECTRONS IN A1203

Since Al<sub>2</sub>O<sub>3</sub> is a good insulator with a band gap of about 9 eV<sup>10</sup> it is desirable to represent the response of its valence electrons on the basis of a quite different model than that used for the conduction band in Al. The model which we have developed for this purpose is related to that employed by Fry 11 in which the ground state wave function of the valence electrons is described in the tight-binding approximation, while excited states are represented by orthogonalized plane waves (OPW). In our use of the model to obtain a dielectric response function we fix the normalization of the OPW excited states by requiring that the sum rule  $\int_{-\infty}^{\infty} d\omega \, \omega_{\text{Im}}[\epsilon(k,\omega)] = 2\pi^2 \text{ne}^2/\text{m}$ , where n is the density of electrons in the valence band. In addition we assume that the solid is uniform and homogeneous. The dielectric response function corresponding to this model solid is convenient and flexible for use, can be fitted to the optical dielectric function in the limit of very long wavelengths (k→0), and describes the single-particle properties of excited electrons. The existence of plasma oscillations emerges naturally as one studies the response of the system to longitudinal electric perturbations. 12

Since a detailed discussion of the insulator model is planned for publication  $^{13}$  we quote here only the results needed for these calculations. The

result required here is the imaginary part of the dielectric response function for the model insulator given by

$$Im[\epsilon(k,\omega)] = \pi ne^{2} \Gamma / h \beta k \Lambda$$
 (20)

where

$$\Gamma = \left\{ \frac{1}{3} \left[ \frac{1}{\{\alpha^2 + (k-p)^2\}^3} - \frac{1}{\{\alpha^2 + (k+p)^2\}^3} \right] - \frac{32\alpha^4}{(\alpha^2 + p^2)^2 (4\alpha^2 + k^2)^2} \left[ \frac{1}{\alpha^2 + (k-p)^2} - \frac{1}{\alpha^2 + (k+p)^2} \right] (21) + \frac{1024 \, pk \, \alpha^8}{(\alpha^2 + p^2)^4 (4\alpha^2 + k^2)^4} \right\} ,$$

$$\Lambda = \left[ \omega_{B} + \frac{h\beta}{m} (k^{2} + \alpha^{2}) \right] \left[ \frac{1}{8\alpha^{5}} - \frac{32\alpha^{3}}{(4\alpha^{2} + k^{2})^{4}} \right]$$
 (22)

and

$$p = [m(\omega - \omega_g)/h\beta]^{1/2} . \qquad (23)$$

Here  $\hbar\omega_{_{\rm B}}$  is the excitation energy of the valence electrons and  $\beta$  and  $\alpha$  are parameters which may be adjusted to make the theory agree with optical dielectric function measurements in the k $\rightarrow$ 0 limit. In the k $\rightarrow$ 0 limit we have

$$\operatorname{Im}\left[\epsilon(0,\omega)\right] = \frac{2^{9}\pi ne^{2}}{3h\beta} \frac{\alpha^{7}}{\left(\omega_{B} + \frac{h\beta}{m}\alpha^{2}\right)} \frac{p^{3}}{\left(\alpha^{2} + p^{2}\right)^{6}} . \tag{24}$$

Given the imaginary part of the dielectric function, Eq. (20), for fixed values of n,  $\beta$ ,  $\omega_0$ , and  $\alpha$  the real part of  $\epsilon(k,\omega)$  may be obtained numerically using the

Kramers-Kronig relation,

$$\operatorname{Re}[\epsilon(k,\omega)] = 1 + \frac{2P}{\pi} \int_{0}^{\infty} d\omega_{0} \frac{\omega_{0} \operatorname{Im}[\epsilon(k,\omega_{0})]}{\omega_{0}^{2} - \omega^{2}} . \tag{25}$$

We fit Eq. (24) to experimental measurements of the imaginary part of the optical dielectric function as measured by Arakawa and Williams  $^{10}$  with 15 valence electrons per Al<sub>2</sub>O<sub>3</sub> unit, density of Al<sub>2</sub>O<sub>3</sub> of 4.05 g/cm<sup>3</sup>,  $\hbar\omega_8 = 9$  eV,  $\beta = \frac{1}{2}$ , and  $\alpha_0 = 0.78$ . Given these values,  $\text{Im}[-\frac{1}{\epsilon(k,\omega)}]$  is calculated from the results of Eq. (20) and Eq. (25). The remaining 9 of the 24 valence electrons in Al<sub>2</sub>O<sub>3</sub> are taken to form a second tight-binding level with  $\hbar\omega_8 = 29$  eV,  $\beta = \frac{1}{2}$ , and  $\alpha_0 = 1.6$ .

#### VI. EXCHANGE CORRECTED DIMFP'S AND FORMULAE FOR THE TABULATIONS

We have included the effect of electron exchange in our calculations in a simple manner based on the form of the Mott formula (non-relativistic Møller formula) for scattering of an incident electron with a free electron. The cross section for finding a scattered electron with energy in the interval W:W+dW is given by 8

$$\frac{d\Phi}{dW} = \frac{\pi e^4}{E} \left[ \frac{1}{W^2} + \frac{1}{(E-W)^2} - \frac{1}{W(E-W)} \right]$$
 (26)

for an incident electron of energy E, except for energies close to W = 0 and W = E. Near W = 0 and W = E the interference term (third term on the right side of Eq. (26)) is effectively zero.

The DIMFP for excitation of an electron from a particular state i may be written in the form

$$\tau_{i}(E, \hbar\omega) = \frac{1}{E} F_{i}(E, \hbar\omega)$$
 (27)

If we assume that the width of the level from which an electron is excited is quite narrow, we obtain from Eq. (27) the DIMFP for production of a secondary electron with energy  $\mathbf{E}_{s}$  as

$$\tau_{i}^{s}(E, E_{s}) = \frac{1}{E} F_{i}(E, E_{i}^{s} + E_{s})$$
 (28)

where  $E_i^{\theta}$  is the binding energy of the i<sup>th</sup> state (a positive quantity). The exchange corrected DIMFP is taken as

$$\tau_{i}^{\text{exc}}(E, h\omega) = \frac{1}{E} \left\{ F_{i}(E, h\omega) + F_{i}(E, E + E_{i}^{8} - h\omega) - \left[ 1 - \sqrt{E_{i}^{8}/E} \right] \left[ F_{i}(E, h\omega) F_{i}(E, E + E_{i}^{8} - h\omega) \right]^{1/2} \right\}. (29)$$

Since  $E\tau_i \approx 1/(\hbar\omega)^2$  for large E and  $\hbar\omega$ , Eq. (29) reduces in this limit to the form given by Eq. (26). The factor  $1-\sqrt{E_i^8/E}$  reduces the contribution of the third term in Eq. (29) as  $E \to E_i^8$ . This form for the exchange corrected DIMFP has been used in our calculations for all the inner shells and for the two valence levels in  $Al_2O_3$  (since our model assumes the width of these levels to be quite narrow).

If we now define the more energetic of the two electrons after collision to be the primary, and account for exchange through Eq. (29), Eq. (3) gives the

the contribution to the inverse mean free path due to excitation of an electron from the ith level as

$$\mu_{i}(E) = \int_{E_{i}^{a}}^{(E+E_{i}^{a})/2} d(\hbar\omega) \tau_{i}^{exc}(E,\hbar\omega) . \qquad (30)$$

Similarly for the stopping power and mean square energy loss per unit path length we have from Eq. (4) and Eq. (5)

$$S(E) = \int_{E_{i}^{8}}^{(E+E_{i}^{\bullet})/2} d(\hbar\omega) \, \hbar\omega \, \tau_{i}^{\text{exc}}(E,\hbar\omega)$$
(31)

and

$$\Omega_{i}^{2}(E) = \int_{E_{i}^{8}}^{(E+E_{i}^{8})/2} d(\hbar\omega)(\hbar\omega)^{2} \tau_{i}^{exc}(E,\hbar\omega) . \qquad (32)$$

The exchange correction for excitation of electrons from the Al conduction band of finite width requires a slightly different form. We take

$$\tau_{ee}^{exc}(\varepsilon, x) = \left\{ \tau_{ee}(\varepsilon, x) + \tau_{s}(\varepsilon, \varepsilon - x) - [1 - (\varepsilon + 1)^{-1/2}] [\tau_{ee}(\varepsilon, x) \tau_{s}(\varepsilon, \varepsilon - x)]^{1/2} \right\}$$
(33)

where  $\tau_{\rm ee}$  and  $\tau_{\rm s}$  are defined in Eq. (12) and Eq. (17). Thus the contributions to the inverse mean free, stopping power, and mean square energy loss per unit path length due to excitation of electrons from the conduction band in Al are calculated from

$$\mu_{ee}(\varepsilon) = \int_{0}^{x/2} dx \, \tau_{ee}^{exc}(\varepsilon, x) , \qquad (34)$$

$$S_{ee}(\varepsilon) = \int_{0}^{x/2} dx \times \tau_{ee}^{exc}(\varepsilon, x) , \qquad (35)$$

and

$$\Omega_{ee}^{2}(\varepsilon) = \int_{0}^{x/2} dx \, x^{2} \, \tau_{ee}^{exc}(\varepsilon, x) \quad . \tag{36}$$

No exchange correction is applied to  $au_{
m pl}$ , Eq. (14). The inverse mean free path contribution due to plasmon excitation is given by

$$\mu_{pl}(\epsilon) = \int_{x_{min}}^{x_{max}} dx \, \tau_{pl}(\epsilon, x) . \tag{37}$$

The contribution to stopping power and mean square energy loss per unit path length are given by

$$S_{pl}(\epsilon) = \int_{x_{min}}^{x_{max}} dx \times \tau_{pl}(\epsilon, x)$$
(38)

and

$$\Omega_{pl}^{2}(\varepsilon) = \int_{x_{min}}^{x_{max}} dx x^{2} \tau_{pl}(\varepsilon, x) . \qquad (39)$$

For the remaining calculations we form the sums

$$S_{\text{exc}}(E) = \sum_{i} S_{i}(E)$$
 (40)

and

$$\Omega_{\rm exc}^2(E) = \sum_{i} \Omega_i^2(E) \tag{41}$$

where the index i includes the terms appropriate for a given solid, including exchange corrections as indicated above. The csda range is calculated from

$$R_{(10)}(E) = \int_{10 \text{ eV}}^{E} dE'/S_{exc}(E')$$
 (42)

corresponding to an electron slowing down in a continuous manner from an energy E to 10 eV. The mean square fluctuation in the csda range based on Eq. (7) is calculated as

$$[\Delta R_{(10)}]_{AV}^{2} = \int_{AV}^{E} dE' \Omega_{exc}^{2}(E') / [S_{exc}(E')]^{3}.$$
 (43)

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#### VIII. ALUMINUM: EXPLANATION OF TABLES

## **GENERAL NOTES**

- 1. Electron energies are measured from the top of the conduction band (Fermi level). The Fermi energy is taken to be  $E_r = 11.6$  eV.
- 2. The density of solid Al is taken to be 2.71 g/cm<sup>3</sup>.
- 3. The computer-printed units are translated as:

4. The numerical print-cut is in the form, e.g.

$$2.8D-01 = 2.8 \times 10^{-1}$$
.

### TABLE 1A - INVERSE MEAN FREE PATH OF ELECTRONS IN ALUMINUM

EL - EL  $\mu_{ee}$  as given by Eq. (34)

PLASMON  $\mu_{pl}$  as given by Eq. (37)

AL(n $\ell$ )  $\mu_{n\ell}^{A1}$  as given by Eq. (30) with Eq. (19)

INVERSE MFP  $\mu$  - total inverse mean free path = sum of individual contributions.

### TABLE 1B - STOPPING POWER OF ALUMINUM FOR ELECTRONS

EL-EL  $S_{ep}$  as given by Eq. (35)

PLASMON

S<sub>pl</sub> as given by Eq. (38)

AL(nt)

 $S_{n\ell}^{A1}$  as given by Eq. (31) with Eq. (19)

STOPPING POWER

S - total stopping power = sum of individual contributions

# TABLE 1C - CSDA RANGES AND STRAGGLING OF ELECTRONS IN ALUMINUM

CSDA RANGE (E TO 10 EV) R<sub>(10)</sub> - the range of an electron in the continuous-slowing-down approximation in going from an energy E to 10 eV, as given by Eq. (42).

MEAN SQUARE ENERGY LOSS

Ω<sup>2</sup> - the mean square fluctuation in the energy loss per unit path length, as given by Eqs. (41), (36), (39) and (5) plus (19).

MEAN SQUARE RANGE FLUCTUATION

 $\left[\Delta R_{(10)}\right]_{AV}^{2}$  - the mean square fluctuation in the range about the mean csda range  $R_{(10)}$ , as given by Eq. (43).

RELATIVE RANGE STRAGGLING  $\{[\Delta R_{(10)}]^{2}_{NV}\}^{1/2}/R_{(10)}$ 

TABLE 1A-IN/ERSE MEAN FREE PATH OF ELECTRONS IN ALUMINUM (DENSITY 2.71G/C43)

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1	INDIVIDUA	EL-EL	00	8230-0	. 6790-0	- 8980-	• 3590-0	98.00-0	6880-0	0-0410	.3580-0	• 720D-0	0-0660 •	0-0564.	0-0000	7680-0	• 22 1D -0	. 6880-0	· 1670-0	0-0099	010001	1200-0	· 603D-0	• 11 8D-0	. 5780-0	0-0496	0550-0	0760-0	1050-0	• 11 80 -0	01000	0940-0	. 0760-0	0280-0	9-00-00	0-0420	01 80 -0	- 01 20-0	. 0060-0	• 98 20 -0	90104	71 80 -0	0-0019	5160-0	· 41 00 -0	0- 0E 0E •	1950-0
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TABLE 1A-INVERSE MEAN FREE PATH OF ELECTRONS IN ALUMINUM (DENSITY 2.716/CH3)

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J dir	AL (28)	.7250-0	.5970-0		2150-0	·0 86D-0	.9780-0	.8740-0	• 7 730-0		A02.	4060-0	.3220-0	.2420-0	-1640-0	0-0690-	010010	01004	8140-0	.7510-0	0-0069.	•632D-0	5750-0	46.80-0	-417D-0	.3680-0	.320D-G	22740-0		-044	.1040-0	.064D-0	0-0070	8840-0	0-0618	-128D-0	0-000	010000	5430-0	4960-0	.4520-0	0-0014	0-06050	2260-0		מפ
NS TO	AL (2P)	.9550-0	.8640-0	0-04040	6140-0	.5380-0	.4650-0	0-0+65.	• 3260-0	C-01070	1380-0	0-00000	.024D-0	.970D-0	9180-0	0-0808.	010044	72001	6860-0	.644D-0	.6040-0	.5660-0	• 52 BD-0	4570-0	4240-0	. 3910-0	• 36 0D-0	0-0625	2710-0	2.2430-03	. 21 50-0	0-0891•		0640-0	0180-0	0-04/6	0-0556	6550-0	9180-0	7830-0	200-0	7180-0	0-0199	0-0-0		40-04
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TABLE 1A-INVERSE MEAN FREE PATH OF ELECTHONS IN ALUMINUM (DENSITY 2,716/CM3)

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TABLE 18-STOPPING POWER OF ALUMINUM (DENSITY 2.716/CM3) FOR FLECTRONS

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AL (2P)	-6500-	0-0690	0-01000	0-0969	0-0669.	.7000-0	.7010-0	.7020-0	0-0104	01000	0-2009	.6970-0	.6950-0	.6920-0	•6860-0	010000	6370-0	6210-0	.6060-0	5910-0	5640-0	.5480-0	•532D-0	.5160-0	0-0664	4640-0	4460-0	.427D-0	8-4090-01	3700-0	.3490-0	3280-0	• 286D-0	.2640-0	•242D-0	0-0022	1740-0	1480-0	.1230-0	•0970-0	0460-0	-0116	-788D-	A 6 0 P
PLASMON	. 5030	-3360-	- 16 3D-	.0550-	.9520-	- 8510-	• 753D-	-6570-	- 30 +U-	1860-	-3000-	-2170-	· 1350-0	.0556-0	0-0879.	• 40 A D = 0	. 7560-0	.6850-0	• 616D-0	.5490-0	4190-0	.356D-0	· 294D-0	• 2330-0	0-1221-	0-0090	.00500	.9500-0	4	. 793D-0	.743D-0	.6940-0	. 5980-0	• 551D-0	• 5050-0	460D-0	3720-0	.3290-0	• 288D-0	-2470-	1670-	-0776	-802D-	AA 10
	7.8750-01	-0690-	-5680-	-4710-	-377E+	· 2860-	-0261·	-0011.	9440	9640-	-7860-	.710D-	•636D-	-5630-	-0556	**************************************	-2910-	-226D-	• 163D-	- 102D-	.983D-	•925D-	-8670-	-8140-	10407	•655D-	-604D-	-5540-	.5050- .4570-	-4100-	-364D-	-318D-	-530D-	-1870-	-1450-	06 30	. 024D-	9840-	-946D-	-0806.	-01/0- 8340-	-6590-	-4980-	O VE
POWER EV/A	2.6700 00		0 00 19	2.5900 0	. 570D 0	. 5510	0 0555.	0 0504	4790	800	.4450 0	.429D 0	4130 0	0 0255	3660 0	0 0055	.3340 0	.319D 0	• 304D o	2457	. 2610 0	.248D 0	• 2340 0	0 0202		1840 0	0 0171.	1580 0		. 1200 0	• 108D 0	0 0360	.0720	0 00 90 •	0480	00220	0 0410	. 003D	0 0266	0.0020	0 00 96	· 908D 0	· 860D 0	CALA.
ENERGY EV	5.100 02		4000	.50D 0	• 60D o	0 002.	0000			2000	·300 0	400 9	.500 0			0 036	0 000·	0 001.	2000		.500 0	0 009	0 002			.1000	.20D 0	90D 0	50000	0 009	0 002.		0 000	0 001.	2000	4000	.50D 0	0 Q09·	0 00 0		0000	.050 0	· 10D 0	150 0

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TABLE 18-STOPPING POWER OF ALUMINUM (DENSITY 2.71G/CM3) FOR ELECTRONS

F.V.	,						4		0.5	0	0	) C	0	0	C	C	20	2 0	0	0	0	C	0	0	20	0	0	0	0	20	0	<b>C</b> (	00	20	0	0	00	<b>&gt;</b> C	0	0	0	0	0	00	<b>&gt;</b> c	, rc
70	15)								Ë,	င္င်င္ပ			5	55	5	0	7 6		9	30	50	90	င္ပ	9	1 6	9	90-	404	ទូ:	100	9.6	-06	500	5	70-	<u>_</u>	2	700	5	9	5	ę	٤	9		ac.
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ECT	рp	INDIVIDUAL	CONTRIBUTIONS	TO THE STOPP	ING POWER IN	UNITS OF EVIA
ENERG RV	POWER EV/A	EL-FL	PLASMON	AL (2P)	AL (25)	AL(15)
7 00 0	6	1 90-0	4D-0	4	7.: 300-02	9.4230-03
4.80D 03	7.085D-01	1.3940-01	1.4690-01	1 1	780-0	5600-0
0 006	9820	0-01/		30.30-	1380-0	6320-0
	0 6	260-0	0-09	2580-	0610-0	6670-0
200	869	0-050	1740-0	2140-	9850-0	7120-0
300 0	2990	85D-0	520-0	1720-	0-0016	
4 00 0	5100	650-0	310-0	1315		
5000	4210-	4 60-0	0-0011			O-US VE
0 009	3380-	280-0	0-016		50000	874D-0
200	25.70-	0-001		0770	200	8980-0
8 0D 0	1780-	0-026		0410	4970-0	9170-0
			0-001	9060	4340-0	9340-0
	0440	1 4 3D - 0	2020-0	8720	3710-0	9490-0
	8840-	1280-0	1860-0	8390-0	3100-0	9620-0
300	815D-	1130-0	170D-0	8070-0	2510-0	9720-0
4000	7480-	0-0660	1550-0	7760-0	1920-0	0-0626
500 0	6820-	0840-0	1400-0	7450-0	1340-0	9840-0
0 00 9	-0819	0710-0	1250-0	0-0512		
7 00 0	.556D-	0570-0	0-0111	0-0000		
8000	4950-	0440	0400		-04 TO	0780-0
0006	4360-	00000		010000	2000	9720-0
000	10875			5760-0	8080-0	9640-0
	36.60-			5500-0	7570-0	9550-0
	10000	A 20-0	0-0810	5240-0	7060-0	9440-0
	1540-0	730D-0	0210-0	4990-0	656D-C	9320-0
2000	1070-0	6220-0	0-0010	4750-0	6070-0	0-0616
009	-0950	515D-C	9880-0	4510	5580-0	00000
7 00 0	0-0200	4120-0	- 880D-C	4280-0	010010	
008	9580-0	3100-0	7740-0		1000	2 - C - C - C - C - C - C - C - C - C -
006	0-0116	2110-0-	10070		11111	8380-0
	0 - 0000	10000	36.00	3 3 4 0 1 - 0	326D-	820D-C
	0 10010	0.000	3710-	3180-0	2820-	-0008
	0-01-6	837D-(	275D-(	2980-0	2390-(	7910-(
00.4	688D-0	7490-(	. 1820-(	.277D-0	-0961	7600-
500	646D-0	.6620-(	.0910-(	.2580-0	1550-	7390-0
009	. 6 05D- 0	.5770-(	00100	238D-0	- OF LI	0.17
7 00 0	.564D-0	.4940-(	.913D-(	2190-0	-0210	10440
900	.524D-0	• 41 3D-(	.8270-(	0-0002.	-0330	
006	• 486D-0	3330-(	. 7420-0	0-0281	10470	10000
000	4470-0	-255D-	0290-		10000	-02 50 ·
001	4100-0	-1790-	-0876		10018	5720-
200	3730-		- 4990-	1001 1001	A 40 00 00 00 00 00 00 00 00 00 00 00 00	5460-
			1440-	0650	-060 B	-5200-
	10101	2000	2690-	0790	-7740-	-494D-
	2425		1950-	0630	-06E 2	.467D-
200	1990	7500-	1220-	40470	-10501	-44 OD-
008	166D-	6830-	.031D-	0310	-0119·	-4130-
006	1330	911D-(	910-	.0160	-638D-	-38.6D-
000	1010-	-553D-	. 9120-	000 O*	•605D-	•359D-
<b>i</b>				•		

RELATIVE RANGE STAAGGLING 2.71G/CH3 IN ALUMINUM (DENSITY SQUARE LUCTUATION A2 MEAN RANGE FL ELECTRONS MEAN SQUARE ENERGY LOSS EV2/A Ь STRAGGL ING RANGE 10EV) **9**¥ CSDA R (E TO IC-CSDA RANGE ELECTRON ENERGY EV TABLE 

	RELATIVE RANGE STAGGLING		1
	FLUCTUATEDA A2		
ECTHONS IN ALUM	QUARE MEA LOSS RANGE		
TRAGGLING OF EL	E MEAN S		
-CSDA RANGE AND S	C SDA RANG (E TO 10E A	88888889999999999999999999999999999999	
TABLE 1C-	SLEC TRON ENERGY EV	\$6.00       \$6.00 <td< td=""><td></td></td<>	

RELATIVE RANGE STA AGGL ING 2.716/CH3) (DENSITA MEAN SQUARE RANGE FLUCTUATION A2 AL UNINUK Z ELECTRONS MEAN SQUARE ENERGY LOSS EV2/A **๛๛๛๚๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛** 14.4 14.6 Ь STRAGGL ING m S A RANGE TO LOEV ON W SO/ TABLE 1C-CSDA RANGE υΞ ELECTRON ENERGY EV 

RELATIVE RANGE STAAGGLING	3.0220-01 3.0160-01 3.0100-01	0-02000	0-0166	9790-0	9730-0	9620-0	9560-0	9440-0	9380-0	92.70-0	0-0916	91 00 - 0	8990-0	• 8930-0	. 8820-0	• 876D-0	8650-0	8600-0	8500-0	8440-0	0-0458	8150-0	8050-0	78.70-0	. 779D-0	. 7710-0	010000	74 70-0	. 7400-0	. 7330-0	7190-0	. 7130-0	0-000	
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AR CAT	444																																	
MEAN SOU RANGE FLUCT A2	4.2850 4.6750 5.0900	52.0	98	8.00 9.00 9.00	- C	377	.000 .000 .000	1 49	307	160	571	.666	.869	.9 77 0 8 0	202	.325	.00	.714	96.6	4.5	000	146	5523	6 <b>4</b> E	.800	-276		.868	.4 56	1074	040	112	.265	
QUARE LOSS	000 000	00	000		00		00	0	00 00	00	00	00	00	00	0	00	20	00	90	00	) c	0	00	00	0	90	30	0	0	20	0	00		)
MEAN S ENERGY EV2	2.996 2.992 2.986	98	97	96	900	9.00		.93	92	6.	90	66	. 89	888	88	9	86	. 86	.86	. 85	֓֞֝֞֜֜֝֞֜֝֓֓֓֓֓֓֓֓֓֓֓֓֓֓֓֓֓֓֓֓֓֓֡֓֓֓֓֓֡֓֓֓֡֓֡֓֡֓֓֡֓֡	8	40.	8	*8	•	9 4	. 84	• 8 •	9	. 8	48	9	
NGE OE V)	222	22	200	250	200	E 0	500	E C	200	£00	i M	E0	E C	n n	E C	500	200	M F	n m	503	7 F	n c	n r	, FO	EO	m r	160	E C	603	500	E S	n i		
CSDA RA (E TO L	6.851D 7.170D 7.496D	• 629D • 1680	5140	. 2250	. 5900	0340	1110	.1510	.2330	2740	3590	04030	4920	5370	6290	.7240	7730	.8210	9210	.9720	800	.2880	39.70	6240	. 7400	0000	000	.2280	.3560	6170	.7510	. 8870	1650	)))
770N 767 V	m m m 000	m m																																
ELEC	1.450 1.550 1.550	୬ ୬	1.700	8	30	•	90		٠,	Q P	) P)	44	S	ບໍ່ຈ	91	٠,	8	80	.0	0-	ď	E,	a R	9	~	00	C	-	'n	.4	S,	0	8	

• 716/C43)	RELATIVE RANGE STAGGLING	6880	6760-	- 6700-	06650	653D-	- 64 80-	. 64 20-	-03 20	42 70 -	-0129	-0919.	-61 1D-	-0909	2000	2000	-583	. 583D-	- 5780-	- 57 30-	- 04 40 ·	56 00-	- 25 60-	- 55 10 -	-04 40	5380-	- 5340-	2000	. 52.20-	- 51 80-	10000	0000	40 AD	-0464	4900-	4820-	4790-	-4750-	44680-	2.4640-01	10001
ITY 2	7	9								:			c					٠,		7										e				c ,	-		2				
DENS	ARE																																								
ALUMINUM (	MEAN SQUI ANGE FLUCTO	310	6 130	0012	0010	0 230	.1360	2530		0000	1660	0906	• 0 50 D		5140	6 790	.8490	.0 250	.2060	0265	7820	9860	0561	0114	0198	0960	.3350	9 3 3 5 0	0960	3620	0210	2040	0108	01110	7 5 2 5	0800	0450	0770	490	1 - 1860	
Z	à	•																																			¢.				:
CI RONS	JARE																																							800	
ING UT ELE	MEAN SOU ENERGY L EV2/A	833	. 833	200	300	.817	.814	100		.802	. 799	. 796	. 793	44	783	. 779	.776	.772	100	900	. 756	. 752	.748	***	1 7 M S	.731	. 727	.718	+14.	000	. 702	600	060	• 686	678	.674	.670	.667	659	2.6560	1
4 73 34 77 .	m S					_												_							_								,								
}	RANG 10E	000	0	<b>&gt;</b> C	0	0	0	0 0	0	0	0	0	00	0	0	0	0	0	> c	0	0	0	00	<b>7</b> C	0	0	00	0	0	00	0	00	0	00	0	0	0	9 0	Ö	86 00	1
1	C SDA (E TO	4.451	. 746	70	200	.361	.519	0 0 0	000	.175	344	. 515		40	.219	.400	. 583		1400	335	. 527	. 722	9		521	. 726	5000	035	.056	000	121	7 € 7 €	188	210	25.0	279	203	0 C	16	98	
;																																		  -  -							
1	S Y	E00																																							
	ELECT ENER EVER	5.100	2,	? •	S	9	- 0	0	9	7	N.	? <	S	9	•	8	2	? -			4	r,	0	8	•	9 -			•	90	-	90	•	-10	10	4	ָּהַ עַ		0	•	

### IX. ALUMINUM OXIDE: EXPLANATION OF TABLES

### **GENERAL NOTES:**

- 1. Electron energies are measured from the bottom of the conduction band.
- 2. The density of solid Al<sub>2</sub>O<sub>3</sub> is taken to be 4.05 g/cm<sup>3</sup>.
- 3. The computer-printed units are +ranslated as:

4. The numerical print-out is in the form, e.g.,  $2.8D-2 \equiv 2.8 \times 10^{-2}$ .

# TABLE 2A - INVERSE MEAN FREE PATH OF ELECTRONS IN $^{12}O_{3}$

VAL (9 EV) 
$$\mu_{(9 \text{ eV})} \text{ as given by Eq. (30)} + (2) + (20) + (25)$$

$$VAL (29 \text{ EV}) \qquad \mu_{(29 \text{ eV})} \text{ as given by Eq. (30)} + (2) + (20) + (25)$$

$$AL(nt) \qquad \mu_{nt}^{A1,A1} = 2^{O_3} \text{ as given by Eq. (30)} + (19)$$

$$O(1S) \qquad \mu_{1s}^{O,A1} = 2^{O_3} \text{ as given by Eq. (30)} + (19)$$

$$INVERSE \text{ MFP} \qquad \mu - \text{total inverse mean free path} = \text{sum of individual contributions.}$$

### TABLE 2B - STOPPING POWER OF ALUMINUM OXIDE FOR ELECTRONS

VAL (9 EV)

 $S_{(9 \text{ eV})}$  as given by Eq. (31) +(2) +(20) +(25)

**VAL (29 EV)** 

 $S_{(29 \text{ eV})}$  as given by Eq. (31) +(2) +(20) +(25)

AL(nt)

 $S_{n,\ell}^{A1,A1} = {}^{O}_{3}$  as given by Eq. (31) +(19)

O(1S)

 $S_{1s}^{O,A1} 2^{O_3}$  as given by Eq. (31)+(19)

STOPPING POWER

S - total stopping power = sum of individual contributions.

# Table 2C - CSDA range and straggling of electrons in $^{\mathrm{Al}}_{2}\mathrm{O}_{3}$

CSDA RANGE (E TO D EV) R<sub>(10)</sub> - the range of an electron in the continuous-slowing-down approximation in going from an energy E to 10 eV, as given by Eq. (42).

MEAN SQUARE ENERGY LOSS

 $\Omega^2$  - the mean square fluctuation in the energy loss per unit path length, as given by Eqs. (41), (32)+(19), and (32)+(2)+(20)+(25).

MEAN SQUARE
RANGE FLUCTUATION

 $\left[ \Delta R_{(10)} \right]_{AV}^{2}$  - the mean square fluctuation in the range about the mean csda range  $R_{(10)}$ , as given by Eq. (43).

RELATIVE RANGE STRAGGLING

 $\{ [\Delta R_{(10)}]_{AV}^{2} \}^{1/2} / R_{(10)}$ 

TABLE 24-INVERSE MEAN FREE PATH OF ELECTRONS IN ALZO3 (DENSITY 4.05G/CM3)

1-V	AL(15)	•	0	•	• •		•	•	•	•	• •		•	•	•			•				•	•	•	• •	• •	•	•	•			•	•			•	•	•	•	• •	• •	•	•	•	•	•
IN UNITS OF	0(15)	•	0	•			•	•	•	•	•		•	•								•	•	•			•	•	• •		•	•	•				•	•	• (			٠	•.	• •		•
HE INVERSE MFP	AL (25)	•	0.0	•	• •	•	•	٠	•	• •	• •	•	•	•	•	• •	•	•	•	•	•	•	•	• •	•	•	٠	•	• •	•	•	•	• •	• •	•	•	•	•	• •	•	•	•	•	• •	•	•
IEUTIONS TO T	AL ( 2P )	•	0.0	•	• •		•	•	•	• •	•	•	•	•	•	•	•	•	•	•	•	•	•	• •	•	•	•	•	• •	٠.	٠	٠	• •	•	•		•	• •	• •		•	•	•	• •	•	•
VIDUAL CONTRI	ш	•	0.0	•	• •	•	•	•	•	• •	•	•	٠	•	• •	•	•	•		•	•	•	•	• •		0	.2610-0		.021D-0	.3650-0	1050-0	0-06/2	2910-0	.8800-0	.5160-0	1940-0	0-01000	0-0045	.0610-0	.7240-0	.3570-0	19560-0	0-0400	• •	0-0660.	.1410-0
101	VAL (9EV)	•	000		• •	•	•	•		3560-0	.1110-0	• 0C2D-0	• 392D-0		50000	8890-0	. 254D-0	.6160-0	• 974D-0	• 33 9D -0	· 698D-0	0490-0	0-0000	0210-0	.3090-0	. 576D-0	• 820 D = 0		- 4400-	. 0050-0	• 474D-0	00000	0-0660 •	. 1420-0	. 1770-0	• 204D-0	10000	3210-	3480-	3700-	-0666	42.30-	46701		4930-	507D-
INVERSE	- 1	0.0	0.0		00	0.0	0.0			.3560-0	.1110-0	.0020-0	3,420-0		5220-0	0-0588	.2540-0	.6160-0	.9740-0	13390-0	0-0369•	40440-0	71.50	.0210-0	.309D-0	.5760-0	0 - 0 - 0 - 0	7910-0	.4660-0	0.0250-0	.561D-0	0.0000	1220-0	.171D-0	.2120-0	.2460-0		3 850-0	.418D-0	.4470-0	820-0		0-089	1.5020-01	0-050	210-0
ELEC TRON	בי שני	3 -000 ·	1.000 00		000	0 000	88			0 000	001.	2000.0		. ני פיני	0009	.7CD 0	.80D O	0 006.	0 000	001.				.60D 0	. 7 CD C	000		200	4 00 0	969		200	4 00 0	.600 0			400	600 0	.80D o	0 000	000		800 0	0 000	2000	0 00 .

TABLE 2A-INVERSE MEAN FREE PATH OF ELECTRONS IN AL203 (DENSITY 4.05G/CM3)

A-1	AL(15)	•	•	•	•		• •	•	•	•	•	•	•	•	•	• •	•	•	•	• •	•	•	•	• •	0	•	• •	•	•	•	• •	•	•	• •		•	٠	• •	•	•	•	• •	•	•
IN UNITS OF	(\$1)0			•	• •			•	•	٠	•	•	•	• :	• •		•	•	• •	•	•	•	• •		0.0	•	•	•	•	•	• •	•				•		•	•	•	•	•		•
αį	AL (25)	•	•	•	•	•	•	•	•	•	•	•	•	2 550	. 9260-0	.745D-0	-3970-	9440-0	00000	.4640-0	.9480-0	4 1 10-0	2610-0	.6450-	6.0030-04	6.5000	-9150-	.9260-	.0270-	1860-	2500-	-3080-	-0505.	4290-	-4610-	4 880	+ 3 I ZU-	-0705	-4800-	• 4 55U-	10000	179	.3530-	9
// i	AL ( 2P )	•	•	4 4 9 0 - 0	6040-0	.5890-0	.2170-0	.9210-0	• 74 ED-0	.6470-0	.5740-0	. 53.20-0	0-0444	0-0-1	.0750-0	·4860-0	0-0268	010010	7680-0	.9560-0	. 10 20-3	4 20-0	. 54 60-0	.6570-0	4 . 74 90-03	.8770-0	.9340-0	.9870-0	0-0250	0-05-0	1120-0	1210-0	1420-0	.1480-0	1520-0	1520-0	01000	1370-0	1280-0	0-0711	0-06-00	.073D-0	.0550-0	0-07 50.
יייייייייייייייייייייייייייייייייייייי	V AL (29EY)	800-0	0-000	200-0	0-050	300-0	530-0	730-0	920-0	0-080	230-0	0-0-0	01000	920-0	370-0	760-0	140-0	2010	320-0	0-016	250-0	0-060	970-0	840-0	1.63.00102 0.00102	380-0	200-0	0-050	670-0	500-0	320-0	140-0	800-0	640-0	500-0	0-0-6	100-0	0-096	030-0	0-050	390-0	24D-0	0-060	020
	VAL (9EV)	•	ם ספר	10-	-06	-08	20-	-06	ב ב ב				30-	70-	20-	5290-	44 70-	4080-	3700-	3330-	10000	2330-	2020-	-1730-	10-0611	. 09 30 -	-0690 •	0470	. 0040-0	. 84 20-0	6490-0	28401	1120-0	9470-0	6480-0	4880-0	. 3450-0	• 208D-0	0.0000	8250-0	. 705D-0	· 2800-0	0-08/6	0-0707
MF P	-	1.6370-01	6660-	790-	-006	7010-	1120-	1022	1000	7440-	7500-	7540-	764D-	7620-	7470-	10000	6530-	-0819	5040-	1000	4040-	453D-	4220-	10445	1390-	-3130-	-2860-	2410-	-2150-	-0261.	-1770-	1370-	-1160-	0000	0660.	0-0360	.0380-0	0-0510.	9090-0	7630-	·627D-0	0-0565	2410-0	
ENERGY	נ	7.600 01	600	200	200	36			0000	60D 0	. BCD C	0 000	0 001.	2002.		2005	.600	2007.	0000		0001	2000 0	0000	5000	. 600	926		000	0 001.	2000		.500	999		0005	0 000.	88		200	\$ 005	0 00 0		000	

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I-4	AL(15)	•	• •	000	•	• •	•	•	•	•	•	•	•	• •	•	•	•	•	•	٠	•	• •	•	•	٠	•	•	> 0	• •		•	•	• •	•	•	٠	•	•	•	•	•	•	• •	•	•	•
INITS OF	0(15)	000		. 8220		2000	9840-0	.5430-0	1950-0	. 874D-0	• 0560 - C		יין כני ער איני יין כני ער איני	7110	.8660-0	0-0910.	1770-0	.3360-0	. 4940-0	2 - GE - G		0 - 09 6 6 •	. 2 3 ED - C	.3760-0	· 31 10-0	• 6430 - 0	0 - 0177		1250-0	.2400-0	0 - 24 80 - 0		6560-0	. 7530-C	6 4 40 - 0		1220-0	2100-0	. 2960-6	3800-0	• 4623-0		7000-0	. 063D-0	.3760-0	• 6210-0
E INVERSE MFP	AL (25)	7.2730-04	.2190-0	0-0261	0 -040	0-0190	0-0510	.9710-0	.9270-0	•883D-0	0.40104	7590-0	7190-0	6800-0	.6420-0	• 6050-0	.5690-0	• 5330- 0	• 4980-0	0 -0000	0 1000	3640-0	.3320-0	· 300D-0	• 269D- 0	0-2662*	0 -080 - 0	1450-0	.1060-0	.0670-0	0-0820	0-0536	.9170-0	.881D-0	0-00-0		7440-0	.7110-0	.6 790-0	.647D-0	0-0919	A 5 5 5 5 5 5 5 5 5 5 5 5 5 5 5 5 5 5 5	.5260-0	.3860-0	.252D-0	0 -0 + 1 1 •
EUTIONS TO TH	AL ( 2P )	W 0 - 00 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0	.9580-0	. 935D-0	010110	86.20-0	.8390-0	. 81 50-0	. 7920-0	0-0011	10000	7010-0	6780-0	6550-0	.6310-0	.6070-0	• 58 0D-0	• 5540-0	0-0876	010000	4560-0	4340-0	4130-0	0-006E •	• 36 8D-0	040745	010706	. 28 70-0	.2640-0	0.2410-0	010017	1730-0	· 150D-0	• 1280-0	010000	0-0190	0390-0	·018D-0	0-0966	0-03/6	0-0466	01000	0920-0	. 7920-0	.6960-0	0-0000
ت ده	V AL (29	1.2670-02	2400-0	270-0	0-01-00	1890-0	1770-0	1650-0	1540-0			1100-0	0-0660	0-0680	0-0610	0-0690	0-0600			0-0220	0130-0	0-08000	·962D-0	.878D-0	0-0967		0-0655	.482D-0	.4070-0	3320-0	1870-0	1160-0	.0400-0	0-0116.	8410-0	.776D-0	.711D-0	.6470-0	•584D-0	0-0226	0-0104	.342D-0	.2840-0	0-0900	0.000-0	
INDI	į	7.1630-02	96 80-0	• E75D-0 • 785D-0	0-02-5	6119-0	. 5280-0	4480-0	0-0000 0-0000	010810	1450-0	.0740-0	· 00500-0	. 537D-0	. 87.20-0	• 6070-0	010010	* 50 4 0 1 0 * 50 4 0 1 0	5640-0	5080-0	4520-0	. 397D-0	3440-0	2910-0	0-00-0	1410-0	. 093D-0	. 0450-0	0-0666	9540-0	. E660-0	. 823D-0	. 78 10-0		6600-0	. 6210-0	. 583D-0	. 5430-0	0-0606	010714	4020-0	3670-0	- 334D-0	0-00-0	8870-0	
INVERSE		9.002D-02 8.0870-02	.776D-0	0-0/00.	• 4 500 - 0	3570-0	.2600-0			0-0866	0.070	.723D-0	0-0119.	.5600-0	• 4 B20 - 0			1840-0	1140-0	.0450-0	.9780-0	.9120-0	. 348D-0	0 - 020 - 0	6620-0	6 030 - 0	.5450-0	4 8 B D - 0	0 - 0 7 6 6	3220-0	.2690-0	2100-0	1650-0	0650	0-0910.	· 968D-0	9200-0	0-04/20	7.440-0	7400-0	6970-0	.6540-0	6120-0	010200	0.260-0	!
ELECTRON FNEDGY	2 2 >	5.100 02 5.200 02	0000	2000	0 009	100	000			2000	300	0 00 4.	.500	009.				0001	2000	300 0	4 00 0	S S S S S S S S S S S S S S S S S S S			006.	0 000	0 031.	0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0		2000	.600 0	88		000	0001	2000			0009	0 002.	0 008	0006		0001	150 0	

p : e'

TABLE 2A-INVERSE MEAN FREE PATH OF ELECTRONS IN AL203 (DENSITY 4.056/CM3)

ELECTRON ENERGY	E STAN STAN STAN STAN STAN STAN STAN STAN		UAL CO	: :	INVERSE MFP	IN UNITS OF	A-1
м >		VAL(9EV)	V AL (29 EV)	AL ( 2P )		0(18)	AL(15)
1.200 03	0970	7590-0	•289D-0	.5190-0	9070-0	. 84 00 - 0	•
300	0-0709	530		3,4370-03	4.8710-04	7,0360-05	0
350	4 750-0	. 426D-0	2-0112-	2830-0	0-0259	0-0012	• 1
	3550-0	. 3280-0	.542D-0	.2110-0	.5150-0	.4680-0	• •
500		0 - 0 - 0 - 0 - 0 - 0 - 0 - 0 - 0 - 0 -	0-0066.	1430-0	.4050-0	. 5410-0	•
550 0	0200-0	0690-0-0	010000	0-02-0	.3110-0	6040-0	0
9600	.9300-0	. 9920-0	9390-0	0-05 56	1310-0	. 00 CO .	08190-
	8370-0	0-0616.	-806p-0	.8940-0	.0470-0	.7450-0	2860-
2000		0-0000	•6790-0	. 838D-0	.9650-0	. 7800-0	.3630-
.800	.5850-0	7210-0	-0300-0 -0-0-0-0	1840-0	6870-0	2:00	. 54 70-
0.000	• 5 0 cp - 0	• 66 20 -0	3310-0	6820-0	7 380-0		-0661
	• 4 30D - 0	• 6050-0	.2250-0	6340-0	6690-0	0-066	3670-
800	0-0000	0-0155	1240-0	. 58 80-0	0-0109.	160-0	. 624 D-
.050	.2330-0	4490-0	0-0-0-0-0-0	010000	9360-0	9360-0	-8370-
86	.1740-0	. 402D-0	.8440-0	4590-0	- 100 - V	56	-0001
	1150-0	· 356D-0	.7580-0	4190-0	3830-0		-0117
		* 31 20-0 212 0	•6780-0	.3810-0	.2970-0	8640-0	4220-
0000	010400	0-0000	.5950-0	.3440-0	.2420-0	370-0	. 5290-
350 0	.9020-0	1910-0	0-0516	0-08060	1 900-0	8060-0	.6310-
0 0000	.8530-0	1530-0	0-0525	010000	0-065	96	- 7290-
0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0	8070-0	.1170-0	3040-0	2086-0	0 10000		. 824D-
	.7610-0	. 0820-0	.237D-0	1770-0	0.00		10000
	6740-0	0400-0	1720-0	.1470-0	.9530-0		0840-
650 0	6330-0	010000	0-0011	11.00-0	0-0016.	6200-0	-1640-
.700 0	. 5960-0	9540-0	010100	0-0060	•8690-0 9369	3880-0	.240D-
6.4 양양	.5580-0	. 9250-0	9340-0	0-0960	0-0004	35 40 1 0 1 0 1 0 1 0 1 0 1 0 1 0 1 0 1 0	-3140-
000	5210-0	· 897D-0	.8790-0	.0110-	7530-0		10000
0006	4 300 - 0	• 869D-0	·8260-0	860-0	. 7170-0	44 80-0	-3060-
950	4170-0	8170-0	0-040-0	0-019	.6810-0	0-07	- 3640-
86	.3840-0	. 7910-0	6750-0	450-0	010444		-02009
26	.3210-0	• 74 3D-0	.5810-0	0-069	5510-0	26.00	7730-
	0 - 0 - 0 - 0 - 0 - 0 - 0 - 0 - 0 - 0 -	0-0860.	•493D-0	260-0	.4910-0	1920-0	. 866D-
400	1530-0	6140-0	4040-0	9 60-0	0-0504	1160-0	- 9480-
2005	· 1 030-0	5760-0	2530-0	100-0	1110-0		-0.120
000	0.0550	539D-0	1810-0	740-0	284D-0	0.000	10700
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98	0-0006.	9-01-0	0-04400	0-080	• 1960-0	. 71 50-0	. 2520-
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8	0-02+8.	3800-0	8670-0	0100	0-071	.5630-0	-0046.
200.0	.8160-0	3530-0	.0120-0	920-0	0.0000		102.00
	7750-0	3260-0	.7890-0	650-0	0-0010	8	440001
30	7100-0	010-010	.708D-0	400-0	750-0	690-0	550
9.009	5750-0		-00000 -00000 -00000	0-09-0	.9420-0	0-0861.	-4760-
			0-0610	0-026	100-0	1300-0	.4890-
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. '	VAL(9EV)	V AL (29EV)	AL ( 2P )	AL ( 23)		AL(15)
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1.6210-02	1.2090-02	2.5240-03	1.3460-03	1.8490-04	£, 9980 - 05	3.5080-06
5940-0	1800-0	.4830-0	3260-0	200-0	9380-0	. 513b-0
2670-0	1680-0	.4420-0	30 ED-0	910-0	.8730-0	.5160-0
	0-0041	0-0504.	2860-0	640-0	. 0120-0	.5180-0
4920-0	2010	0-0006	010000	380-0	4 5 40 1 Q	0-06-15
4 70D-0	0940-0	2000	01016	1001		
4470-0	0440	26.00	0 1 0 1	0 - 0 - 0	36.5	7-0010
4200-0	0-0-0-0	000000	1070			
0 - O : O 4		010101		0 - 0 5 5	010000	3-0206.
		0-0061	0-000	0-0/1	• 4 B B - 0	4940-0
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	0-0c10 •	0-040-0	0-0661	130-0	• 3800 - 0	.4770-0
0-0055	• 001D-0	0-0401	1340-0	520-0	.3320-0	.4670-0
3210-0	• 86 7 D -0	.0750-0	1190-0	320-0	2850-0	4570-C
0-050 F	. 7330-0	.049D-0	1050-0	120- C	2.100-0	4450-0
29-0-0	.6010-0	.0220-0	0-0160	0-050	2001	
27.0-0	0-0576	0.0400				
2010			01000	0-002		0-0024.
		0-0116	0-000	4300-0	0-001	. 408D-0
0-0252	. 2280-0	0-09460	5 20-0	4390-0	.0660-0	395D-0
0-0/22	• II 00-0	220-0	4 0D-0	422D-0	. 0260-0	3010-0
2120-0	0-0966 •	0-066	0280-0	4050-0		7666
1970-0	- 884 D-0	760-0	200			
1820-0	7760-0			000000000000000000000000000000000000000	010144	0-0565
		01010	0-0500	3720-0	0-0606 •	33300-0
	0-01/0	320-0	0-08 66	3570-0	.878-0	. 325D-C
0-0001	• 567D-0	110~0	023D-0	3420-0	9360-0	3100-0
1410-0	.4660-0	0-016	.7160-0	1270-0	0100	0.000
1280-0	. 3690-0	7.0	100	1000		
1160-0	0-0F 4C				01001	0-0002
	010000	01010	-0010	2980-0	O-OFF	.2650-0
	0-0001	340-0	-0114	2840-0	2000	. 2500-0
0-0160	0-0060 •	130-0	-3140-	2710-	6.00-0	2380-0
0-0510	00100 •	950-0	-219D-	2570-	0100	22000
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0-0950	010000			1000		3-0007
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		201	100 17	-0617	200-0	1 74 0-0
	0-0400	260-0	•8580 <del>-</del>	2070-	230-0	.1590-0
0-0000	· 2820-0	100-0	-0577	1950-	0-056	. 144D-0
0-0410.	. 5080-0	940-0	-6890-	1940	24.70	2000
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0-0360-	1440-0	190-0	-2960-	129D-	280-0	0-0530
.5660-0	0.0220	0-050	.2220-0	1001	200	
.4710-0	0-0600 •	0-016	1490-0	2000		
.3830-0	0-01 30	770-0	01000			
2980-0	0.00 P			10000	2 6 6 6 6 6 6 6 6 6 6 6 6 6 6 6 6 6 6 6	0-0000
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AL (15) EVIA 9 FOR ELECTRONS UNITS (51)0 z POWER 4.05G/CM3) AL (25) STOPPING **CDENSITY** THE AL(2P) CONTRIBUTIONS TO ALUMINUM OXIDE VAL (29EV) INDIVIDUAL 6 TABLE 28-STOPPING POWER VAL(9EV) STOPPING POWER EV/A 0000000000 0C 20000000000000 GL ECTRON ENERGY EV 

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EV/A AL(15)	•	•	•	•	•	•	•	•	•	•	• •	•	•	•	•	•	•	•	•	•	•	•	• •	• •		•	٠	•	•	•		•	•	•	•	•			•	•	•	•	•	•	•	•		•
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STOPPING POWER	6	6	5	96	96	5				100	70-	50-	30-	010	-066	970-	3950-	-0£6£	3920-	10000	10000	1000	3820-	3800-	378D-	3760-	100	2000	1024	3630-	1.3590-01	3540-	3500-	10047	-025	3330-	3290-	3250-	32 10-	-0715	1000	1000	100 P	-086	940-	760-	580-	4004
710NS TO THE AL(2P)	-380D-	-3910-	-3665	-0000	10000			4000	3070-	3930-	-3066-	-3050-	-380D-	-3740-	-367D-	-3609-	• 3520-	-33300-	1000	1000	27.0	2560-	.2420-	-2290-	-2130-	-19801•	-1820-	1000	1430	1250-	7.1070-01	-0880-	-0000	1000	-0600	-9880-	-0996	-9450-	-9230-		85,701	-05180	-8130-	-0061.	-768D-	. 655	.543	• 4 32
DUAL CONTRIBU	-07470-	-6750-	-0040-	-0256	10101	32.0-	2500-	-184D-	-116D-	-0490-	.982D-	-0116.	•852D-	-0882-	- 7250-	-0630-	-0100	-00400	10000	100 25	3050	-248D+	· 1920-	-136D-	-0810-	-0280-	10000	A700-	-8190-	-768D-	7.7180-01	-66690-	- 0 C C C C	5240-	-4770-	-430D-	•384D-	-3380-	-0565	10400	1610-	-118D-	-0920-	• 034D-	. 9920-	-793D-	-0000-	0150
INDIVI VAL (9EV)	7610 0	. 7260 0			2000	5650	5360 0	.5070	0 0674.	4510 0	.4240 0	.3980 0	.3730 0	.34ED 0	0 0000			0 0000	י מטוכי	1880	1670	0 0441.	.1270 0	• 108D 0	• 0880 ¢			0 0910	9980	• 981D 0	96.	0480		9010	.8850	.870D O	•856D 0	04100		0000	7850	.773D 0	.7600 0	470 0	340 0	740		
STOPPING POWER EV/A		0 00/5		4550	4180 0	.3810 0	. 345D 0	.310D 0	.275D 0	.241D 0	. 208D o	0 0921	0 0641.	0 0511.					9350	0 0206	.880D O	.853D 0	.8270 0	0 0108	0 002.	0 0922	7020 0	.6780	• 6550 O	•632D 0	0 060		5410	.5190	.4970 0	4760 0		0 0 0 0 0 0 0	O OF OF	0 0575	.3530 0	• 333D 0	.3140 0	.295D 0	0 00/2.			
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8888 0000 AL (15) EVZA 6 ELECTRONS IN UNITS FOR POWER 4.05G/CH3) STOPPING AL(2S) THE 22.22.20 22.22.20 23.22. (DENS 10 6570-AL ( 2P) CONTRIBUTIONS OX I DE ALUMI NUM AL (29EV INDIVIDUAL **P** POWER 9890-01 9630-01 9530-01 9120-01 8880-01 8640-01 AL (9EV) 28-STOPPING MUNNUM MUNUM MUNUM MUNUM MUNUM 11.1860 11.1860 11.1860 11.1860 11.181 STOPPING POWER EV/A TABLE ELECTRON ENERGY EV 

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TABLE 2C-CSDA RANGE AND STRAGGLING OF ELECTRONS IN AL203 (DENSITY 4.05G/CH3)

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	-400 01 1-1060 02 1-9150 02 1-1770 04 9-8100- -600 01 1-1130 02 2-0030 02 1-1760 04 9-7490- -800 01 1-1130 02 2-0030 02 1-1760 04 9-7490- -000 02 1-1170 02 2-0450 02 1-1760 04 9-7190-	-400 01   1-1060 02   1-9150 02   1-1770 04   9-7790-0   9-7790-0   1-1100 02   1-1770 04   9-7790-0   9-7790-	1.106D 02 1.999D 02 1.999D 02 1.17D 02 2.003D 02 1.17D 04 9.77D 04 9.779D-0	100   100	9 6 6 0 0 1 1 1 1 0 6 0 0 2 1 1 9 9 9 0 0 1 1 1 1 1 1 1 1 1 1 1 1 1

RELATIVE RANGE STRAGGLING	5.6150-01	.513D-0	34630-0	4140-0	010000	2 720-0	.2260-0	• 1 820-0	0-0350-0	0520-0	.01110-	0-0016	0-000	.853D-0	.6150-0	7.420-0	.7060-0	.6710-0	6370-0	5700-0	.538D-0	€ 5060-0 • 5060-0	4 450-0	4150-0	.3850-0	3560-0	3000-0	.2730-0	.2460-0	•220D-0	1 640-0	1440-0	.1200-0	0.020-0	0-0000	7 190-0	.640D-0	.5680-0	CHOCK
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C SDA RANGE (E TO 10EV) A	2.018D 02 2.041D 02	0 0880	.1110 0	0 0321.	0 0581	.208D O	.232D 0	.257D 0	0 0202	3330	.3580 0	• 384D • • 600	4360 0	.462D 0	• 488D 0	5420 0	.5690 0	• 596D 0	• 6230 0 • 6500 0	.678D 0	.706D 0	.7340 0	0 0067	. 8190	.8470 0	0 0 0 0 0	9350	.9640	0 0E66.	0230		1130 0	.1440 0	0 0662		.7890	.961D 0	1370 0	2 0411
ELECTRON ENERGY EV	5.800 02 5.900 02	0001	.200 0	9300	5000	• 600 0	0 002	0 000	0000	0 00 1	•200 0 200 0	0000	.500	·60D 0	0 00 0	0 006	0 000	0001		0 000	. 50D 0		8000	0 006	0 000		3000	. 40D 0	9200		8000	0 006.	0 000		2005	.200 0	250 0	0 005	0 0 0 0 0

05G/CM3)	RELATIVE RANGE STRAGGLING	3370-	2 100- 1 740-	4 ~ D !	0000	9705-	9360-	6910-	. 8 5 CD	8 320-	. 8 1 30- . 8 0 4 0- . 7 9 50-	7870	7645-	7370-	- 0 30 - 0 4 0 4 0 4 0 4 0 4 0 4 0 4 0 4 0 4	6750-	6580-	6420-	6280-	6 1 40-	. 6 020- . 5 9 60-
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ABLE	RON 64	mm mm	mmr	m7m	mmn 000	<u> </u>	776	m m r	M M	m m r	) E E	m m r ο τ/ ο	mm:	200	E 0 0	n m r	300	E 60	n m :	nn 00	E 0
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RELATIVE RANGE STRAGGLING	8890 8780 8730 8730	8 4 8 9 9 9 9 9 9 9 9 9 9 9 9 9 9 9 9 9	00000000000000000000000000000000000000	00000000000000000000000000000000000000	10000 10000 10000 10000 10000
MEAN SOUARE RANGE FLUCTUATION A?	0520 3600 6810 0160	215000000000000000000000000000000000000	244448898 244498898 244889898 2448989898989898989898989898989898989898	444446 4	3090
MEAN SQUARE ENERGY LOSS EV2/A	4010 4010 6010 6010 6010 6010 6010 6010			444444 24444 24444 24444 24444 24444 24444 24444 24444 24	662000
CSDA RANGE (E TO 10EV) A	1450 9240 0150 1080 2010	00000000000000000000000000000000000000	24444444444444444444444444444444444444	5.5770 6.1990 6.1990 6.1990 6.1990 6.1990 6.1990 6.1990 7.2890 7.2890 7.6670 6.390 6.300 6.300 6.300 6.300 6.300 6.300 6.300 6.300 6.300 6.300 6.300 6.300 6.300 6.300	3790 5250 6720 0
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